

Copper–indium ordering in $RECu_6In_6$ ($RE = Y, Ce, Pr, Nd, Gd, Tb, Dy$)

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Abstract

The rare earth metal–copper–indides $RECu_6In_6$ ($RE = Y, Ce, Pr, Nd, Gd, Tb, Dy$) were synthesized from the elements by arc-melting. Well-crystallized samples were obtained by slowly cooling the melted buttons from 1320 to 670 K in sealed silica tubes in a muffle furnace. They were investigated by X-ray diffraction on powders and single crystals: $ThMn_{12}$ type, space group $I4/mmm$, $Z = 2$, $a = 916.3(2)$, $c = 535.8(2)$ pm, $wR_2 = 0.063$, 216 F^2 values, 15 variables for YCu_6In_6 , $a = 926.5(4)$, $c = 543.5(3)$ pm, $wR_2 = 0.064$, 314 F^2 values, 15 variables for $CeCu_6In_6$, $a = 925.7(4)$, $c = 540.1(3)$ pm, $wR_2 = 0.075$, 219 F^2 values, 15 variables for $PrCu_6In_6$, $a = 923.1(4)$, $c = 540.3(3)$ pm, $wR_2 = 0.071$, 218 F^2 values, 15 variables for $NdCu_6In_6$, $a = 917.7(4)$, $c = 540.2(3)$ pm, $wR_2 = 0.076$, 207 F^2 values, 15 variables for $GdCu_6In_6$, $a = 917.0(5)$, $c = 540.5(4)$ pm, $wR_2 = 0.062$, 215 F^2 values, 15 variables for $TbCu_6In_6$, $a = 915.2(8)$, $c = 540.7(7)$ pm, $wR_2 = 0.108$, 218 F^2 values, 15 variables for $DyCu_6In_6$. The structures have been refined with a split position (50% Cu + 50% In) for the $8j$ site. They can be explained by a tetragonal body-centered packing of CN 20 polyhedra (10Cu + 10In) around the rare earth atoms. The ordering models of the copper and indium atoms and the limitations/resolution of X-ray diffraction for this topic are discussed.

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1. Introduction

Except scandium, the rare earth metal (RE)–copper–indium systems have been thoroughly investigated in the last 30 years and the isothermal sections of the respective phase diagrams at 670 and 870 K (depending on the rare earth metal content) have been constructed. An overview on the crystal chemistry and physical properties of the many $RE_xCu_yIn_z$ phases is given in a recent review article [1].

So far more than 100 ternary compounds have been reported for these ternary systems, but only few of them have completely been characterized on the basis of single crystal diffractometer data. The series of $RECu_{5.1}In_{6.9}$

($RE = Y, Ce, Pr, Nd, Sm, Gd-Lu$) [2,3] indides crystallizes with the tetragonal $ThMn_{12}$ type structure [4], space group $I4/mmm$. So far, several other series with compositions RET_4X_8 (T = late transition metal; X = element of the 3rd main group), RET_5X_7 , RET_6X_6 , and $RET_{10}X_2$ have been reported [5]. With respect to composition, several of them have no possibility for $T-X$ ordering. Keeping the high tetragonal space group symmetry, a fully ordered structure is possible only for the composition RET_4X_8 . When lowering the space group symmetry from $I4/mmm$ via a *translationengleiche* symmetry reduction of index 2 (t_2) to $Immm$, an ordered composition RET_6X_6 is possible through splitting of the $8j$ subcell site into the fourfold sites $4f$ and $4h$. This kind of ordering is realized for the series of $REFe_6Ga_6$ gallides [6] and has first been observed for $ScFe_6Ga_6$, $ZrFe_6Ga_6$, and $HfFe_6Ga_6$ [7].

Refinement of the $ErCu_{5.1}In_{6.9}$ structure on the basis of X-ray powder data revealed the following site occupancies:

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the rare earth atoms on the 2*a* site, In on 8*i*, and the mixed occupied sites 8*f* (Cu_{0.85}In_{0.15}) and 8*j* (Cu_{0.422}In_{0.578}) [3]. Since the composition of these indides is close to $RE\text{Cu}_6\text{In}_6$, we were interested in the precise Cu/In ordering. The crystal growth and single crystal structure refinements of $RE\text{Cu}_6\text{In}_6$ ($RE = \text{Y, Ce, Pr, Nd, Gd, Tb, Dy}$) indides are reported herein.

2. Experimental

2.1. Synthesis

Starting materials for the preparation of the $RE\text{Cu}_6\text{In}_6$ ($RE = \text{Y, Ce, Pr, Nd, Gd, Tb, Dy}$) indides were ingots of the rare earth elements (Johnson Matthey, Chempur or Kelpin), copper wire (Johnson Matthey, $\varnothing 1\text{ mm}$), and indium tear drops (Merck), all with stated purities better than 99.9%. The rare earth metal pieces were first arc-melted [8] to small buttons under ca. 600 mbar argon atmosphere. The argon was purified over titanium sponge (900 K), silica gel, and molecular sieves. The rare earth metal buttons were then mixed with pieces of the copper wire and the indium tear drops in the 1:6:6 atomic ratios and the mixtures were reacted in the arc-furnace and remelted three times to ensure homogeneity. The weight losses after the arc-melting procedures were always less than 0.5%. The light gray polycrystalline samples are brittle and stable in air over months. Finely ground powders are dark gray and single crystals exhibit metallic luster. For crystal growth, the arc-melted buttons were sealed in evacuated silica tubes and first rapidly heated at 1320 K in a muffle furnace. The temperature was then lowered at a rate of 5 K/h to 1020 K and then at a rate of 15 K/h to 670 K. The polycrystalline samples are stable in moist air over months.

2.2. X-ray powder data

The $RE\text{Cu}_6\text{In}_6$ samples were characterized through Guinier powder patterns using $\text{CuK}\alpha_1$ radiation and α -quartz ($a = 491.30$, $c = 540.46\text{ pm}$) as an internal standard. The Guinier camera was equipped with an imaging plate system (Fujifilm, BAS-1800). The tetragonal lattice parameters (Table 1) were refined by least-squares calcula-

tions using the WinXPow software supplied by Stoe. To ensure correct indexing, the experimental patterns were compared with calculated ones [9], taking the atomic sites obtained from the structure refinements.

2.3. Single crystal X-ray diffraction

Irregularly shaped crystals of $RE\text{Cu}_6\text{In}_6$ ($RE = \text{Y, Ce, Pr, Nd, Gd, Tb, Dy}$) were selected from the crushed annealed samples. These crystals were glued to small quartz fibres using bees wax and first checked by Laue photographs on a Buerger camera, equipped with the same Fujifilm, BAS-1800 imaging plate technique. Good quality crystal of CeCu_6In_6 , PrCu_6In_6 , GdCu_6In_6 , and TbCu_6In_6 were used for the intensity data collection on a Stoe IPDS II diffractometer (graphite monochromatized $\text{MoK}\alpha$ radiation; oscillation mode). Numerical absorption corrections were applied to the data sets. The YCu_6In_6 , NdCu_6In_6 , and DyCu_6In_6 crystals were measured at room temperature by use of a four-circle diffractometer (CAD4) with graphite monochromatized $\text{MoK}\alpha$ radiation and a scintillation counter with pulse height discrimination. Scans were taken in the $\omega/2\theta$ mode. Empirical absorption corrections were applied on the basis of Ψ -scan data, accompanied by spherical absorption corrections. Relevant crystallographic data for the data collections and evaluations are listed in Tables 2 and 3.

2.4. Scanning electron microscopy

The single crystals investigated on the diffractometer were analyzed using a LEICA 420 I scanning electron microscope with CeO_2 , the rare earth trifluorides, Cu, and InAs as standards. No impurity elements heavier than sodium were observed. The compositions determined by EDX ($10 \pm 2\text{ at\% Y}$: $45 \pm 2\text{ at\% Cu}$: $45 \pm 2\text{ at\% In}$, $8 \pm 2\text{ at\% Ce}$: $45 \pm 2\text{ at\% Cu}$: $47 \pm 2\text{ at\% In}$, $8 \pm 2\text{ at\% Pr}$: $44 \pm 2\text{ at\% Cu}$: $48 \pm 2\text{ at\% In}$, $8 \pm 2\text{ at\% Nd}$: $45 \pm 2\text{ at\% Cu}$: $47 \pm 2\text{ at\% In}$, $8 \pm 2\text{ at\% Gd}$: $47 \pm 2\text{ at\% Cu}$: $45 \pm 2\text{ at\% In}$, $7 \pm 2\text{ at\% Tb}$: $46 \pm 2\text{ at\% Cu}$: $47 \pm 2\text{ at\% In}$, $8 \pm 2\text{ at\% Dy}$: $47 \pm 2\text{ at\% Cu}$: $45 \pm 2\text{ at\% In}$) are close to the ideal composition, i.e. 7.6:46.2:46.2. The standard uncertainties account for the analyses at various points (between 3 and 10 points per sample).

3. Results and discussion

3.1. Structure refinements

Careful analyses of the IDPS data sets revealed high Laue symmetry and body-centered tetragonal cells. There were no further systematic extinctions, leading to the space groups $I4/mmm$, $I4mm$, $I4_2m$, $I4m2$ and $I422$, of which the centrosymmetric group was found to be correct during structure refinement. The atomic positions of $\text{ErCu}_{5.1}\text{In}_{6.9}$ [3] were taken as starting values and the seven structures were refined with anisotropic displacement parameters for

Table 1

Lattice parameters (Guinier powder data) of the ternary indium compounds $RE\text{Cu}_6\text{In}_6$ ($RE = \text{Y, Ce, Pr, Nd, Gd, Tb, Dy}$)

Compound	<i>a</i> (pm)	<i>c</i> (pm)	<i>V</i> (nm ³)
YCu_6In_6	916.3(2)	535.8(2)	0.4499
CeCu_6In_6	926.5(4)	543.5(3)	0.4665
PrCu_6In_6	925.7(4)	540.1(3)	0.4628
NdCu_6In_6	923.1(4)	540.3(3)	0.4604
GdCu_6In_6	917.7(4)	540.2(3)	0.4549
TbCu_6In_6	917.0(5)	540.5(4)	0.4545
DyCu_6In_6	915.2(8)	540.7(7)	0.4529

Table 2

Crystal data and structure refinement for YCu₆In₆, CeCu₆In₆, PrCu₆In₆, and NdCu₆In₆, space group *I4/mmm*, *Z* = 2

Empirical formula	YCu ₆ In ₆	CeCu ₆ In ₆	PrCu ₆ In ₆	NdCu ₆ In ₆
Molar mass	1159.07 g/mol	1210.28 g/mol	1211.07 g/mol	1214.40 g/mol
Unit cell dimensions	Table 1	Table 1	Table 1	Table 1
Calculated density	8.56 g/cm ³	8.62 g/cm ³	8.69 g/cm ³	8.76 g/cm ³
Crystal size	30 × 40 × 190 μm ³	10 × 30 × 100 μm ³	40 × 70 × 120 μm ³	20 × 40 × 170 μm ³
Detector distance	–	60 mm	60 mm	–
Exposure time	–	5 min	5 min	–
ω range; increment	–	0–180°; 1.0°	0–180°; 1.0°	–
Integr. param. A, B, EMS	–	13.0; 3.0; 0.014	13.5; 3.5; 0.012	–
Transm. ratio (max/min)	1.57	2.21	1.96	7.03
Absorption coefficient	35.2 mm ^{−1}	32.6 mm ^{−1}	33.2 mm ^{−1}	33.7 mm ^{−1}
<i>F</i> (0 0 0)	1014	1052	1054	1056
θ range	3–30°	3–35°	3–30°	3–30°
Range in <i>hkl</i>	±12, ±12, +7	±14, ±14, ±8	±12, ±12, ±7	±12, ±12, ±7
Total no. reflections	1435	3513	2413	2634
Independent reflections	216 (<i>R</i> _{int} = 0.034)	314 (<i>R</i> _{int} = 0.038)	219 (<i>R</i> _{int} = 0.062)	218 (<i>R</i> _{int} = 0.081)
Reflections with <i>I</i> > 2σ(<i>I</i>)	204 (<i>R</i> _σ = 0.015)	236 (<i>R</i> _σ = 0.043)	207 (<i>R</i> _σ = 0.025)	213 (<i>R</i> _σ = 0.027)
Data/parameters	216/15	314/15	219/15	218/15
Goodness-of-fit on <i>F</i> ²	1.179	0.966	1.154	1.279
Final <i>R</i> indices [<i>I</i> > 2σ(<i>I</i>)]	<i>R</i> ₁ = 0.028 <i>wR</i> ₂ = 0.062 <i>R</i> 1 = 0.029 <i>wR</i> ₂ = 0.063	<i>R</i> ₁ = 0.028 <i>wR</i> ₂ = 0.062 <i>R</i> 1 = 0.042 <i>wR</i> ₂ = 0.064	<i>R</i> ₁ = 0.034 <i>wR</i> ₂ = 0.073 <i>R</i> 1 = 0.036 <i>wR</i> ₂ = 0.075	<i>R</i> ₁ = 0.029 <i>wR</i> ₂ = 0.071 <i>R</i> 1 = 0.029 <i>wR</i> ₂ = 0.071
Extinction coefficient	0.0103(6)	0.0088(5)	0.023(2)	0.0063(5)
Largest diff. peak and hole	3.51/−2.92 e/Å ³	4.77/−2.47 e/Å ³	3.09/−2.62 e/Å ³	3.14/−2.93 e/Å ³

Table 3

Crystal data and structure refinement for GdCu₆In₆, TbCu₆In₆, and DyCu₆In₆, space group *I4/mmm*, *Z* = 2

Empirical formula	GdCu ₆ In ₆	TbCu ₆ In ₆	DyCu ₆ In ₆
Molar mass	1227.41 g/mol	1229.08 g/mol	1232.66 g/mol
Unit cell dimensions	Table 1	Table 1	Table 1
Calculated density	8.96 g/cm ³	8.98 g/cm ³	9.04 g/cm ³
Crystal size	10 × 40 × 50 μm ³	20 × 40 × 190 μm ³	20 × 40 × 60 μm ³
Detector distance	60 mm	60 mm	–
Exposure time	25 min	8 min	–
ω range; increment	0–180°; 1.0°	0–180°; 1.0°	–
Integr. param. A, B, EMS	14.0; 4.0; 0.038	14.0; 4.0; 0.022	–
Transm. ratio (max/min)	1.94	4.68	2.89
Absorption coefficient	35.7 mm ^{−1}	36.2 mm ^{−1}	36.8 mm ^{−1}
<i>F</i> (0 0 0)	1064	1066	1068
θ range	2–30°	3–30°	3–30°
Range in <i>hkl</i>	±12, ±12, ±7	±12, ±12, ±7	±12, ±12, ±7
Total no. reflections	1691	2303	2634
Independent reflections	207 (<i>R</i> _{int} = 0.066)	215 (<i>R</i> _{int} = 0.031)	218 (<i>R</i> _{int} = 0.065)
Reflections with <i>I</i> > 2σ(<i>I</i>)	175 (<i>R</i> _σ = 0.033)	212 (<i>R</i> _σ = 0.013)	205 (<i>R</i> _σ = 0.028)
Data/parameters	207/15	215/15	218/15
Goodness-of-fit on <i>F</i> ²	1.174	1.172	1.149
Final <i>R</i> indices [<i>I</i> > 2σ(<i>I</i>)]	<i>R</i> ₁ = 0.035 <i>wR</i> ₂ = 0.073 <i>R</i> 1 = 0.044 <i>wR</i> ₂ = 0.076	<i>R</i> ₁ = 0.030 <i>wR</i> ₂ = 0.062 <i>R</i> 1 = 0.031 <i>wR</i> ₂ = 0.062	<i>R</i> ₁ = 0.045 <i>wR</i> ₂ = 0.104 <i>R</i> 1 = 0.048 <i>wR</i> ₂ = 0.108
Extinction coefficient	0.0020(4)	0.0096(6)	0.013(1)
Largest diff. peak and hole	4.14/−3.30 e/Å ³	4.61/−3.66 e/Å ³	4.82/−3.92 e/Å ³

all atoms with SHELXL-97 (full-matrix least-squares on *F*_o²) [10]. Refinement of the occupancy parameters revealed full occupancy of the 8*f* and 8*i* sites with copper and

indium, respectively. In contrast, the 8*j* sites revealed an occupancy by ca. 50% copper and 50% indium and extremely anisotropic displacement parameters *U*₁₁

(ca. five times larger than U_{22}). Since copper (117 pm) and indium (150 pm) have different covalent radii [11], the extreme anisotropic displacement is a clear hint of copper–indium ordering. Since we observed half occupancy by copper and indium on the $8j$ sites, a subgroup with a splitting into two fourfold sites could solve the problem. This splitting is only possible in space group $Immm$ [12], a *translationengleiche* subgroup of index 2 (t_2) of $I4/mmm$. Consequently we refined all data sets in this lower symmetry space group, also accounting for twinning

Table 4
Atomic coordinates and isotropic displacement parameters (pm^2) of $RE\text{Cu}_6\text{In}_6$ ($RE = \text{Y, Ce, Pr, Nd, Gd, Tb, Dy}$)

Atom	Wyckoff site	x	y	z	U_{eq}
YCu₆In₆					
Y	2a	0	0	0	121(5)
Cu	8f	1/4	1/4	1/4	121(4)
In	8i	0.34049(9)	0	0	125(3)
M1 ^a	8j	0.2776(3)	1/2	0	113(4)
M2 ^a	8j	0.3170(3)	1/2	0	113
CeCu₆In₆					
Ce	2a	0	0	0	42(3)
Cu	8f	1/4	1/4	1/4	85(3)
In	8i	0.34226(8)	0	0	97(2)
M1 ^a	8j	0.2779(2)	1/2	0	62(3)
M2 ^a	8j	0.3105(2)	1/2	0	62
PrCu₆In₆					
Pr	2a	0	0	0	46(4)
Cu	8f	1/4	1/4	1/4	83(4)
In	8i	0.34213(9)	0	0	94(4)
M1 ^a	8j	0.2777(3)	1/2	0	63(4)
M2 ^a	8j	0.3115(3)	1/2	0	63
NdCu₆In₆					
Nd	2a	0	0	0	66(4)
Cu	8f	1/4	1/4	1/4	100(4)
In	8i	0.3417(1)	0	0	110(4)
M1 ^a	8j	0.2783(3)	1/2	0	89(4)
M2 ^a	8j	0.3119(3)	1/2	0	89
GdCu₆In₆					
Gd	2a	0	0	0	83(5)
Cu	8f	1/4	1/4	1/4	117(5)
In	8i	0.3406(1)	0	0	122(4)
M1 ^a	8j	0.2779(4)	1/2	0	93(5)
M2 ^a	8j	0.3165(4)	1/2	0	93
TbCu₆In₆					
Tb	2a	0	0	0	91(4)
Cu	8f	1/4	1/4	1/4	117(4)
In	8i	0.3401(1)	0	0	121(4)
M1 ^a	8j	0.2783(3)	1/2	0	99(4)
M2 ^a	8j	0.3177(3)	1/2	0	99
DyCu₆In₆					
Dy	2a	0	0	0	90(6)
Cu	8f	1/4	1/4	1/4	117(6)
In	8i	0.3399(2)	0	0	114(5)
M1 ^a	8j	0.2778(4)	1/2	0	89(6)
M2 ^a	8j	0.3192(4)	1/2	0	89

U_{eq} is defined as one third of the trace of the orthogonalized U_{ij} tensor.

^aThe M1 and M2 sites are statistically occupied with 25% Cu and 25% In. These sites were refined with equal displacement parameters.

(010, 100, 00–1) due to the *translationengleiche* symmetry reduction. All data sets revealed batch scale factors of 0.5 within one standard deviation and fully occupied sites, but the anisotropic displacements were not well resolved. The copper site still had a high U_{22} parameter, while the U_{22} and U_{33} parameters of indium refined to zero. Furthermore, we observed non-negligible positive and negative residual densities in the final Fourier maps. This readily indicated too small domain sizes, although the samples were very slowly cooled down to 670 K. Thus, we observe no amplitude but phase addition. Consequently, we refined the structures with a split model with Cu/In statistics (50/50) and equal (constrained) isotropic displacement parameters. The refinements then smoothly converged to the residuals listed in Tables 2 and 3. The crystal chemical consequences and reasons for this behavior are discussed in the following section. Final difference Fourier syntheses revealed no significant residual peaks. The atomic parameters and interatomic distances are listed in Tables 4

Table 5
Interatomic distances (pm), calculated with the powder lattice parameters of CeCu_6In_6

Ce	4	In1	317.1
	8	M2	323.5
	8	M1	340.9
	8	Cu1	354.6
In1	4	Cu1	281.8
	1	In1	292.3
	2	M1	293.7
	2	M1	296.1
	2	M2	306.4
	1	Ce	317.1
	2	M2	322.7
	4	In1	341.4
Cu1	4	M1	269.8
	2	Cu1	271.8
	4	M2	274.3
	4	In1	281.8
	2	Ce	354.6
	2		
M1	1	M2	[30.2]
	4	Cu1	269.8
	2	M2	270.5
	2	M1	291.0
	2	In1	293.7
	2	In1	296.1
	2	Ce	340.9
	1	M2	381.3
M2	1	M1	[30.2]
	2	M2	248.2
	2	M1	270.5
	4	Cu1	274.3
	2	In1	306.4
	2	In1	322.7
	2	Ce	323.5
	1	M2	351.1
	1	M1	381.3

All distances within the first coordination spheres are listed. Standard deviations are ≤ 0.2 pm. Note that the M positions have only partial occupancy (see Table 4).

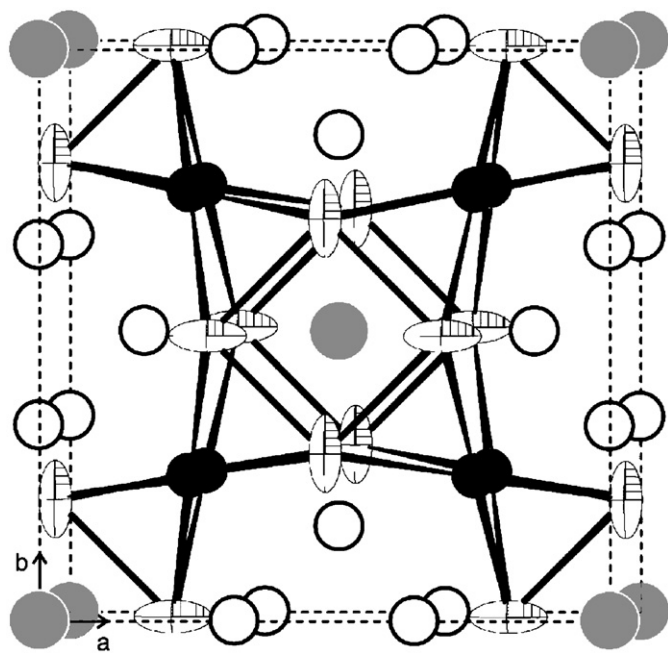


Fig. 1. The crystal structure of CeCu_6In_6 . Cerium, copper, and indium atoms are drawn as medium gray, filled, and open circles, respectively. The copper–indium network around the cerium atoms is emphasized. The displacement ellipsoids of the mixed occupied Cu/In site are drawn at the 99% probability level.

and 5. Further data on the structure refinement are available.¹

3.2. Crystal chemistry

The indides RECu_6In_6 ($\text{RE} = \text{Y}, \text{Ce}, \text{Pr}, \text{Nd}, \text{Gd}, \text{Tb}, \text{Dy}$) were structurally characterized on the basis of single crystal diffractometer data. Although we used very slow cooling rates down to 670 K, no long-range ordering of the copper and indium atoms was observed. The crystals had too small domain size, and consequently, the split model with Cu/In statistics was the best description for the investigated crystals. Nevertheless, we can assume a high degree of short-range order.

As a representative we present a view of the CeCu_6In_6 unit cell with the parameters obtained from the refinement without split positions in Fig. 1. Each cerium atom has coordination number (CN) 20 and the whole structure can be understood as a tetragonal body-centered packing of CN 20 polyhedra. In Fig. 2 the $xy0$ plane of this refinement result is shown. The elongated cigar-like shape of these $\text{Cu}_{0.5}\text{In}_{0.5}$ positions can be resolved in different ways. Keeping the high tetragonal symmetry, we can describe the structure only with Cu/In split positions with a

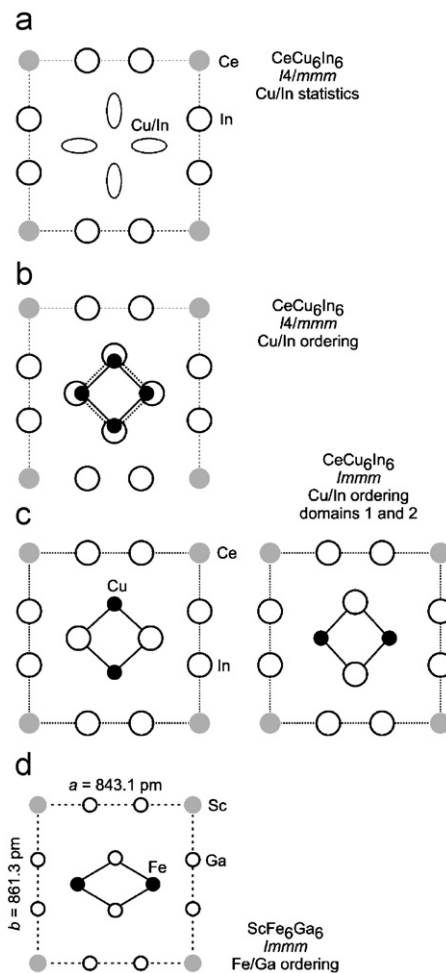


Fig. 2. View of the copper–indium network at $z = 0$ in the structure of CeCu_6In_6 . Cerium, copper, and indium atoms are drawn as medium gray, filled, and open circles, respectively. The statistical model, a split model, and the different ordered models are presented together with the structure of ScFe_6Ga_6 [7]. For details see text.

statistical occupancy. This statistical model can be ordered in two ways: (i) an inner square uniquely occupied by copper and an outer one uniquely occupied by indium atoms or (ii) two Cu_2In_2 rhombs that are rotated by 90° . Since the split refinements clearly revealed full statistics, we can safely assume that the high symmetry model very well reproduces the overlap of the Cu_2In_2 rhombs rotated by 90° . This is a consequence of the twinning observed due to the *translationengleiche* symmetry reduction (the corresponding group–subgroup scheme is shown in Fig. 3). Refinement of the structures with the $\text{Cu}_{0.5}\text{In}_{0.5}$ cigars in all cases resulted in too small In–In distances in the range from 260 to 265 pm (depending on the size of the rare earth element). The split model presented here for the RECu_6In_6 indides has also been observed for the structure of $\text{CaCu}_{6.06}\text{In}_{5.94}$ [13]. A possible reason for the splitting is to allow the formation of suitable bond lengths between the copper and indium atoms. Similar split models also occurred for the BaHg_{11} -related phases $\text{Y}_3\text{TaNi}_{6+x}\text{Al}_{26}$ [14] and $\text{M}_3\text{Au}_{6+x}\text{Al}_{26}\text{Ti}$ ($\text{M} = \text{Ca}, \text{Sr}, \text{Yb}$) [15].

¹Details may be obtained from: Fachinformationszentrum Karlsruhe, D-76344 Eggenstein-Leopoldshafen (Germany), by quoting the Registry Nos. CSD–391410 (YCu_6In_6), CSD–391411 (CeCu_6In_6), CSD–391412 (PrCu_6In_6), CSD–391413 (NdCu_6In_6), CSD–391419 (GdCu_6In_6), CSD–391420 (TbCu_6In_6), CSD–391421 (DyCu_6In_6).

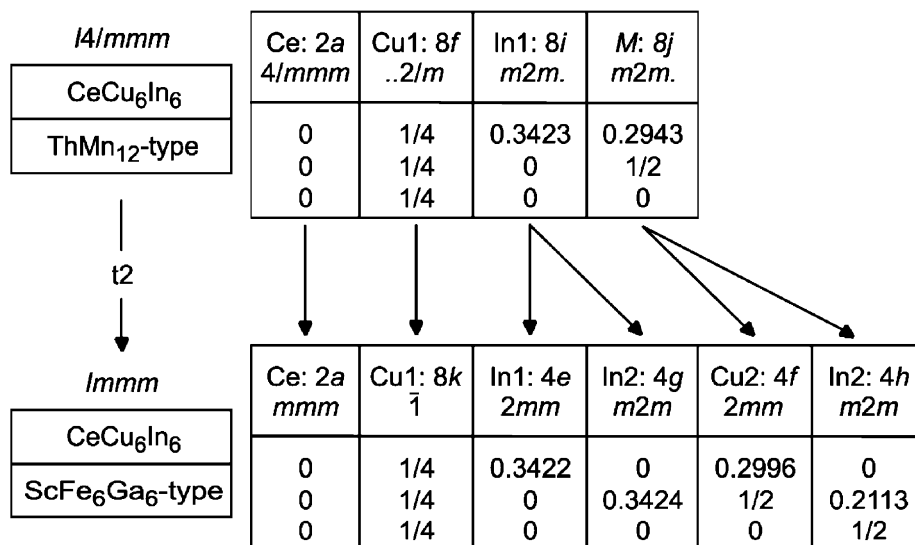


Fig. 3. Group-subgroup relation between the structures of ThMn₁₂ and ScFe₆Ga₆ in the *Bärnighausen* formalism [20,21]. The index of the *translationengleiche* symmetry reduction and the evolution of the atomic parameters are also given. For details see text.

The model with the fully ordered Cu₂In₂ rhombs corresponds to the ScFe₆Ga₆ type [7], space group *Immm*; however, there is a severe difference between this gallide and the indides reported herein. The *a* (843.1 pm) and *b* (861.3 pm) lattice parameters of ScFe₆Ga₆ [7] clearly show a strong orthorhombic distortion, similar to the series of the REFe₆Ga₆ gallides with the heavier rare earth elements [6]. For the RECu₆In₆ indides, neither the arc-melted nor the annealed samples gave indication for splitting of reflections in the powder patterns. Thus, we have no hint for an orthorhombic distortion. As already shown for the series of REFe₆Ga₆ gallides, formation of the superstructure is a function of the size of the rare earth element and temperature as well. For the indides the [Cu₆In₆] network is much larger than the [Fe₆Ga₆] network and this is most likely the reason that we still observe the tetragonal cell parameters for the indides.

From the refinement in space group *Immm* (ordered model) we can calculate the Cu–Cu, Cu–In, and In–In distances (as an example for the cerium compound) within the [Cu₆In₆] network. The Cu–Cu distances of 272 pm are slightly longer than in *fcc* copper (256 pm) [16]. The various Cu–In distances range from 270 to 282 pm, close to the sum of the covalent radii of 267 pm [11]. The shortest In–In distances (291–313 pm) are all shorter than in tetragonal body-centered indium (4 × 325 and 8 × 338 pm) [16]. Similar short In–In distances occur in most indium rich RE_xT_yIn_z compounds, e.g. the series RE₄Pt₁₀In₂₁ [17] or HT-GdNiIn₂ [18]. We can thus assume significant Cu–In and In–In besides weaker Cu–Cu bonding in these RECu₆In₆ indides.

The indides REAg₆In₆ (RE = La, Ce, Pr, Nd, Sm, Eu) [19] show similar ThMn₁₂-related unit cells. Precise single crystal studies with AgK α radiation (since silver and indium differ by only two electrons) are in progress in order to determine the silver–indium ordering.

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